Fundamental interaction studies with nuclear β decay

21st Colloque GANIL

Strasbourg, Sept 2019



The search for 'New Physics'



New Physics experimental searches...

- Energy frontier → LHC, ...
 - Intensity frontier → Nuclear physics, muon, ...
- Cosmic frontier → Planck, ...

NEW PHYSICS : a new theory that completes the SM and solves (at least some of) the current puzzles.







New Physics searches with β decays





- **Specific model;** Beg et al. (1977), Barbieri et al. (1985), Marciano & Sirlin (1987), Hagiwara et al. (1995), Kurylov & Ramsey-Musolf (2002), Marciano (2007), Bauman et al. (2012), ...
- Something more model-indep? EFTs!

Implications for New Physics?

200

Channe

300



[Hardy & Towner'15]

New Physics searches with β decays $fill = \frac{1}{2} \int \frac{1}{2} \int$



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Chann

Competitive probes?

n

- Other low-E searches
- High-E (LHC!!)



 $[V_{ud} = 0.97416(21)!!!]$

[Hardy & Towner'15]

New Physics searches with β decays $f_{\overline{v}}$ f_{\overline



- Specific model; Beg et al. (1977), Barbieri et al. (1985), Marciano & Sirlin (1987), Hagiwara et al. (1995), Kurylov & Ramsey-Musolf (2002), Marciano (2007), Bauman et al. (2012), ...
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Competitive probes?

n

- Other low-E searches
- High-E (LHC!!)

Very active field!

["Recent" review: MGA, O. Naviliat Cuncic, N. Severijns, Prog. Part. Nucl. Phys. 104 (2019) 165-223] ... outdated a few months afterwards!

 $V_{ud} = 0.97416(21)!!!$

[Hardy & Towner'15]

What's an EFT?



Effective Field Theory = Fields + Symmetries - nuclei, e, v - Lorentz - QED - hadrons, e, v - q, u, d, l, e - SU(2) x U(1) - W, Ζ, γ, g - Flavour sym? - B, L;



Not assumption independent!

What's an EFT?



Comparing experiments

- How to compare different nuclear beta decays?
 - Effective Lagrangian at the hadron level! \rightarrow

$$\begin{aligned} -\mathcal{L}_{n \to p e^- \bar{\nu}_e} &= \bar{p} \ n \ (C_S \bar{e} \nu_e - C'_S \bar{e} \gamma_5 \nu_e) \\ &+ \bar{p} \gamma^\mu n \left(C_V \bar{e} \gamma_\mu \nu_e - C'_V \bar{e} \gamma_\mu \gamma_5 \nu_e \right) \\ &+ \frac{1}{2} \bar{p} \sigma^{\mu\nu} n \left(C_T \bar{e} \sigma_{\mu\nu} \nu_e - C'_T \bar{e} \sigma_{\mu\nu} \gamma_5 \nu_e \right) \\ &- \bar{p} \gamma^\mu \gamma_5 n \left(C_A \bar{e} \gamma_\mu \gamma_5 \nu_e - C'_A \bar{e} \gamma_\mu \nu_e \right) \\ &+ \bar{p} \gamma_5 n \ \left(C_P \bar{e} \gamma_5 \nu_e - C'_P \bar{e} \nu_e \right) + \text{h.c.} \end{aligned}$$

• How to compare with e.g. pion decays?

Effective Lagrangian at the quark level! \rightarrow

$$\mathcal{L}_{d \to u \ell^- \bar{\nu}_{\ell}} = -\frac{4G_F V_{ij}}{\sqrt{2}} \left[\bar{\ell}_L \gamma_\mu \nu \cdot \bar{u} \gamma^\mu d_L + \sum_{\rho \delta \Gamma} \epsilon^{\Gamma}_{\rho \delta} \bar{\ell}_{\rho} \Gamma \nu \cdot \bar{u} \Gamma d_{\delta} \right]$$

 $C_i \sim FF \ge \varepsilon_i$

• How to compare with LHC experiments?

Effective Lagrangian at the quark level at the EW scale! \rightarrow

$$\mathcal{L}_{eff.} = \mathcal{L}_{SM} + rac{1}{\Lambda^2} \sum lpha_i \mathcal{O}_i$$







[Lee & Yang'1956]

Hadrons: $n \rightarrow p e^{-} \overline{\nu}$



$$\begin{aligned} -\mathcal{L}_{n \to p e^- \bar{\nu}_e} &= C_V \left(\bar{p} \gamma^\mu n \, + \, \frac{C_A}{C_V} \, \bar{p} \gamma^\mu \gamma_5 n \right) \, \times \, \bar{e} \gamma_\mu \left(1 - \gamma_5 \right) \nu_e \\ &+ C_S \, \bar{p} \, n \, \times \, \bar{e} \left(1 - \gamma_5 \right) \nu_e \, + \, \frac{1}{2} \, C_T \, \bar{p} \sigma^{\mu\nu} n \, \times \, \bar{e} \sigma_{\mu\nu} \left(1 - \gamma_5 \right) \nu_e \\ &- C_P \, \bar{p} \gamma_5 n \, \times \, \bar{e} \left(1 - \gamma_5 \right) \nu_e + \text{h.c.} \end{aligned}$$

+ terms with RH neutrinos

$$\begin{aligned} -\mathcal{L}_{n \to p e^- \bar{\nu}_e} &= \underbrace{C_V \left(\bar{p} \gamma^\mu n + \frac{C_A}{C_V} \bar{p} \gamma^\mu \gamma_5 n \right) \times \bar{e} \gamma_\mu (1 - \gamma_5) \nu_e}_{+ C_S \bar{p} n \times \bar{e} (1 - \gamma_5) \nu_e + \frac{1}{2} C_T \bar{p} \sigma^{\mu\nu} n \times \bar{e} \sigma_{\mu\nu} (1 - \gamma_5) \nu_e}_{- C_P \bar{p} \gamma_5 n \times \bar{e} (1 - \gamma_5) \nu_e + \text{h.c.}}_{+ \text{ terms with RH neutrinos}} \end{aligned}$$

$$-\mathcal{L}_{n \to p e^- \bar{\nu}_e} = \left(C_V \left(\bar{p} \gamma^\mu n + \frac{C_A}{C_V} \bar{p} \gamma^\mu \gamma_5 n \right) \times \bar{e} \gamma_\mu (1 - \gamma_5) \nu_e \right) \\ + C_S \bar{p} n \times \bar{e} (1 - \gamma_5) \nu_e + \frac{1}{2} C_T \bar{p} \sigma^{\mu\nu} n \times \bar{e} \sigma_{\mu\nu} (1 - \gamma_5) \nu_e \\ - C_P \bar{p} \gamma_5 n \times \bar{e} (1 - \gamma_5) \nu_e + \text{h.c.} \right)$$

Linear approx: $SM + small + (small)^2$ (Or simply no v_R : $C_i = C_i'$)

$$-\mathcal{L}_{n \to pe^- \bar{\nu}_e} = \begin{pmatrix} C_V \left(\bar{p} \gamma^{\mu} n + \frac{C_A}{C_V} \bar{p} \gamma^{\mu} \gamma_5 n \right) \times \bar{e} \gamma_{\mu} (1 - \gamma_5) \nu_e \\ + C_S \bar{p} n \times \bar{e} (1 - \gamma_5) \nu_e + \frac{1}{2} C_T \bar{p} \sigma^{\mu\nu} n \times \bar{e} \sigma_{\mu\nu} (1 - \gamma_5) \nu_e \\ - \frac{C_P \bar{p} \gamma_5 n \times \bar{e} (1 - \gamma_5) \nu_e + \text{h.e.}}{\langle + \text{ terms with RH neutrinos}} & \text{``since the nucleons are treated nonrelativistically, the pseudoscalar couplings are omitted''} \\ \frac{Linear approx:}{SM + small + (small)^2} & \text{``Wrong reason... } C_P = 348(11) \epsilon_P \\ (\text{Or simply no } \nu_R: C_i = C_i') & \text{``SM terms} \end{pmatrix}$$

Real reason: the bounds on ε_p from pion decays are much stronger!!!

$$|\mathcal{A}(\pi \to \ell \nu)|^2 \sim m_\ell^2 \left(1 + \frac{M_{QCD}}{m_\ell} \epsilon_P\right)^2$$

$$\begin{aligned} -\mathcal{L}_{n \to p e^- \bar{\nu}_e} &= C_V \left(\bar{p} \gamma^\mu n + \frac{C_A}{C_V} \bar{p} \gamma^\mu \gamma_5 n \right) \times \bar{e} \gamma_\mu (1 - \gamma_5) \nu_e \\ &+ C_S \bar{p} n \times \bar{e} (1 - \gamma_5) \nu_e + \frac{1}{2} C_T \bar{p} \sigma^{\mu\nu} n \times \bar{e} \sigma_{\mu\nu} (1 - \gamma_5) \nu_e + \text{h.c.} \end{aligned}$$

Hadronic EFT



Hadronic EFT



Hadronic EFT



[+ CPV effects]





Observed direction of beta emission in

mirror-reversed

Direction of electron

flow through the solenoid coils



e

Observed direction of beta emission in Direction of electron nirror-reversed flow through the lencid coils



✓ Indirect effect in the Ft-values & neutron lifetime:



$$\delta au_n, \delta \mathcal{F}t ~\sim~ -b ~\langle rac{m_e}{E_e}
angle$$





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$$\delta au_n, \delta \mathcal{F}t ~\sim~ -b ~\langle rac{m_e}{E_e}
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Current data

Precision:

0(0.01 - 1)% !!







[Hardy-Towner'2015]

Current data (+ TH!!)

[MGA, O. Naviliat Cuncic, N. Severijns, Prog. Part. Nucl. Phys. 104 (2019) 165-223]



Precision:

0(0.01 - 1)% !!

[Hardy-Towner'2015]

Th: QED + Isospin symmetry breaking corrections

$$\mathcal{F}t_i \equiv ft_i (1 + \delta'_R)(1 + \delta_{NS} - \delta_C)$$

Current data (+ TH!!)

[MGA, O. Naviliat Cuncic, N. Severijns, Prog. Part. Nucl. Phys. 104 (2019) 165-223]



Nuclear structure-dep. corrections ? [Seng, Gorchtein, & Ramsey-Musolf, PRD100 (2019)] [Gorchtein, PRL123 (2019)] Precision:

0(0.01 - 1)% !!







Driven by Ft's, Tn, An!

[MGA, O. Naviliat Cuncic, N. Severijns, Prog. Part. Nucl. Phys. 104 (2019)]

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$\begin{pmatrix} C_A/C_V \\ C_S/C_V \\ C_T/C_A \end{pmatrix} = \begin{pmatrix} -1.2728(17) \\ 0.0014(12) \\ 0.0020(22) \end{pmatrix} \text{ with } \rho = \begin{pmatrix} 0.08 & 1.00 \\ 0.94 & 0.08 & 1.00 \\ -0.32 & 0.85 & -0.31 & 1.00 \end{pmatrix}$	$\left(\begin{array}{c} C_V \\ C_A/C_V \\ C_S/C_V \\ C_T/C_A \end{array}\right) =$	$\left(\begin{array}{c} 0.98595(34)G_F/\sqrt{2}\\ -1.2728(17)\\ 0.0014(12)\\ 0.0020(22) \end{array}\right)$) with $\rho = \begin{pmatrix} 1.00 \\ 0.08 & 1.00 \\ 0.94 & 0.08 \\ -0.32 & 0.85 \end{pmatrix}$	$1.00 \\ -0.31 1.00$
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• One can trivially calculate the precision needed in any other observable to compete:

Example:

$$b_{GT} = f(C_i) \rightarrow \delta b_{GT} = 0.004$$

Reachable! (NSCL, UW-Seattle?, GANIL?)!



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[MGA, O. Naviliat Cuncic, N. Severijns, Prog. Part. Nucl. Phys. 104 (2019)]













Interlude I: mirror beta decays?

- β transitions between isobaric analog states in T = 1/2 isospin doublets; $\rightarrow |M_F|^2 = 1$ and nonzero F/GT mixing ratio (neutron!).
- Many per-mil level determinations of the Ft values! (Exp + Th) [Severijns et al, PRC78 (2008); Hayen & Severijns, 1906.09870; etc.]
- But mixing ratios are unknown
 - \rightarrow another observable (asymmetries!) needed;

$^{A}_{J}$ Decay	$\mathcal{F}t$ [sec]	asymmetry
$^{19}_{1/2}\mathrm{Ne} \to \mathrm{F}$	1721.44(92) [10]	$A_{\beta,0} = -0.0391(14)$
$^{21}_{3/2}$ Na \rightarrow Ne	4071(4) [11]	$\tilde{a}_{\beta\nu} = 0.5502(60)$
$^{35}_{3/2}\mathrm{Ar} ightarrow \mathrm{Cl}$	5688.6(7.2)	$\tilde{A}_{\beta} = 0.430(22)$
$^{37}_{3/2}\mathrm{K} \to \mathrm{Ar}$	4605.4(8.2) [12]	\tilde{A}_{β} =-0.5707(19) [13]

NOTE: LPCTrap analysis ongoing for Ne-19 & Ar-35

SM analysis: [Naviliat-Cuncic & Severijns, PRL102 (2009)]
 V_{ud} can be extracted with 0.1% precision!
 Although (*currently!*) not competitive, it's a nontrivial crosscheck;

- What about BSM? [Falkowski, MGA & Naviliat-Cuncic, work in progress]
 - In the absence of RH neutrinos, the situation is much like in the SM;
 - Once RH neutrinos are introduced that's not the case.

Interlude II: CP violation?

- If the EFT coefficients are complex, CP-violating effects appear;
- CP violation of "standard" origin is way too small in β decays;
- Beta decay data (D & R correlations) \rightarrow

 $Im(C_A/C_V) = -0.00034(59),$ $Im(C_S/C_V) = -0.007(30),$ $Im(C_T/C_A) = 0.0004(33),$ [MGA, O. Naviliat Cuncic, N. Severijns, Prog. Part. Nucl. Phys. 104 (2019) 165-223]

• Improvements also expected:

Coefficient	Precision goal	Experiment (Laboratory)		Comments
D	$O(10^{-4})$ [418]	MORA (GANIL/JYFL) [418]		In preparation (²³ Mg)
R D, R	O(10 ⁻³) [427] O(0.1)% [399]	MTV (TRIUMF) [427–429] BRAND (ILL) [399,400]		Data taking ongoing (⁸ Li) Proposal
		,	ANR funde	ed!
]	[E. Liéna	rd's talk, Thu 11:10]

Quarks (low-E): $d \rightarrow u e^{-} \overline{\nu}$





[Lifetime shift] $|V_{ud}|^2 + |V_{us}|^2 + |V_{ub}|^2 \neq 1$

$$C_V \sim g_V G_F^{\mu} V_{ud} (1 + \text{NP}) (1 + \text{RC})$$

$$C_A / C_V \sim -g_A / g_V (1 - 2\epsilon_R)$$

$$C_S \sim g_S \epsilon_S$$

$$C_T \sim q_T \epsilon_T$$

$$\tilde{V}_{ud} \equiv V_{ud} \left(1 + \epsilon_L + \epsilon_R\right) \left(1 - \frac{\delta G_F}{G_F}\right)$$

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[Lifetime shift]
$$|V_{ud}|^2 + |V_{us}|^2 + |V_{ub}|^2 \neq 1$$

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 $\begin{array}{rcl}
\tilde{V}_{ud} \\
C_V &\sim & g_V G_F^{\mu} [V_{ud} \left(1 + \mathrm{NP}\right) \left(1 + \mathrm{RC}\right) \\
C_A/C_V &\sim & -g_A/g_V \left(1 - 2\epsilon_R\right) \\
C_S &\sim & g_S \epsilon_S \\
C_T &\sim & g_T \epsilon_T
\end{array}$ [Lifetime shift] $|V_{ud}|^2 + |V_{us}|^2 + |V_{ub}|^2 \neq 1$

Inner RC:

2.361(38)% [Marciano-Sirlin, PRL96 (2006)] 2.467(22)% [Seng et al., PRL121 (2018)] 2.426(32)% [Czarnecki et al., 1907.06737]



$$\tilde{V}_{ud} \equiv V_{ud} \left(1 + \epsilon_L + \epsilon_R\right) \left(1 - \frac{\delta G_F}{G_F}\right)$$

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BUT...

- PNDME'18: $q_A = 1.218(39)$
 - "We argue that our error estimate is more realistic"
- FLAG average: $g_A = 1.251(33)$









[Bhattacharya, Cirigliano, Cohen, Filipuzzi, MGA, Graesser, Gupta, Lin, PRD85 (2012)]





Using these RC + charges, the Ci bounds translate into...



Using these RC + charges, the Ci bounds translate into...



Using these RC + charges, the Ci bounds translate into...





[MGA, O. Naviliat Cuncic, N. Severijns, Prog. Part. Nucl. Phys. 104 (2019)]







Benchmark numbers (from ongoing / planned experiments):

$$\delta \tau_n = 0.1 s$$

 $\tilde{A}_n, a_n, \tilde{a}_F, a_{GT} \text{ at } 0.1\%$
 $b_{GT} = 0.001$

Coefficient	Precision goal	Experiment (Laboratory)	Comments	"Future"
τ _n	$\begin{array}{c} 1.0s; \ 0.1s[210]\\ 1.0s; \ 0.3s[214]\\ 0.2s[215]\\ 0.3s[201]\\ 0.1s[222]\\ \lesssim 0.1s[223]\\ 0.5s[225]\\ 1.0s; \ 0.2s[188] \end{array}$	BL2, BL3 (NIST) [210] LiNA (J-PARC) [211,214] Gravitrap (ILL) [203,215] Ezhov (ILL) [201] PENeLOPE (Munich) [222] UCNτ (LANL) [188,189,223,224] HOPE (ILL) [188,225,226] τ SPECT (Mainz) [188,227]	In preparation; two phases In preparation; two phases Apparatus being upgraded Under construction Being developed Ongoing Proof of principle Ref. [226] Taking data; two phases	<i>π</i> →evγ
β-spectrum β-spectrum b _{GT} b _n	<pre>\$\mathcal{O}(0.01) [256]\$</pre>	Supercond. spectr. (Madison) [256] Si-det. spectr. (Saclay) [253,254] Calorimetry (NSCL) [115,260] miniBETA (Krakow–Leuven) [263–265,270] UCNA-Nab-Leuven (LANL) [271,272,276] UCNA (LANL) [390] PERKEO III (ILL) [295] Nab (LANL) [188,289,357,358] PERC (Munich) [291,292]	Shape factor Eq. (51). Ongoing Shape factor Eq. (51). Ongoing Analysis ongoing (⁶ He, ²⁰ F) Being commissioned Analysis ongoing (⁴⁵ Ca) Ongoing with A_n data Possible with A_n data In preparation Planned	$0^+ \rightarrow 0^+, \tau_n, A_n$
a _F a a _{GT}	0.1% [306] 0.1% [343] 0.1% [79] not stated $\mathcal{O}(0.1)$ % [315]	TRINAT (TRIUMF) [306,310] TAMUTRAP (TA&M) [343] WISArD (ISOLDE) [79,177] Ne-MOT (SARAF) [311,312] ⁶ He-MOT (Seattle) [313,315]	Planned (³⁸ K) Superallowed β p emitters In preparation (³² Ar β p decay) In preparation (¹⁸ Ne, ¹⁹ Ne, ²³ Ne) Ongoing (⁶ He)	
a _{mirror} ã _n a _n	not stated 0.5% [182] 0.5% [273] 1.0% [350] 1.0 - 1.5% [351] 0.15% [188,358]	EIBT (Weizmann Inst.) [316–318] LPCTrap (GANIL) [182,321,323,324] NSL-Trap (Notre Dame) [273,344,345] aCORN (NIST) [350,352–354] aSPECT (ILL) [228,229,351] Nab (LANL) [188,289,357,358]	In preparation (°He) Analysis ongoing (⁶ He, ³⁵ Ar) Planned (¹¹ C, ¹³ N, ¹⁵ O, ¹⁷ F) Data taking ongoing Analysis being finalized In preparation	1 0.000 0.001 0.002 ϵ_T
Α _n Ã _{mirror}	0.14% [391] 0.18% [295] Ø(0.1)% [78]	UCNA (LANL) [390] PERKEO III (ILL) [295] TRINAT (TRIUMF) [78]	Data taking planned Analysis ongoing Planned	mbers / planned experiments):
<i>B</i> _n	0.01% [397]	UCNB (LANL) [397]	Planned	S
$ \begin{array}{l} \tilde{A}_n \left(a_n, \tilde{B}_n, \ldots \right) \\ \tilde{A}_n \left(a_n, \tilde{B}_n, \ldots \right) \end{array} $	0.05% [291] <∅(0.1)% [399]	PERC (Munich) [291,292] BRAND (ILL/ESS) [399,400]	In preparation Proposed	\tilde{a}_{F} acre at 0.1%
D R D, R	$\mathcal{O}(10^{-4})$ [418] $\mathcal{O}(10^{-3})$ [427] $\mathcal{O}(0.1)$ % [399]	MORA (GANIL/JYFL) [418] MTV (TRIUMF) [427–429] BRAND (ILL) [399,400]	In preparation (²³ Mg) Data taking ongoing (⁸ Li) Proposal	001

[MGA, O. Naviliat Cuncic, N. Severijns, Prog. Part. Nucl. Phys. 104 (2019) 165-223]



Quarks, W, Z, ...



Matching with high-E EFT



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NP searches in β decays

 $\frac{d\,\vec{\epsilon}(\mu)}{d\log\mu} = \left(\frac{\alpha(\mu)}{2\pi}\gamma_{\rm ew} + \frac{\alpha_s(\mu)}{2\pi}\gamma_s\right)\,\vec{\epsilon}(\mu),$



$$\begin{split} \frac{\delta G_F}{G_F} &= 2 \ [\hat{\alpha}_{\varphi l}^{(3)}]_{11+22} - [\hat{\alpha}_{ll}^{(1)}]_{1221} - 2[\hat{\alpha}_{ll}^{(3)}]_{1122-\frac{1}{2}(1221)} ,\\ V_{1j} \cdot \epsilon_L^{j\ell} &= 2 \ V_{1j} \ \left[\hat{\alpha}_{\varphi l}^{(3)} \right]_{\ell\ell} + 2 \ \left[V \hat{\alpha}_{\varphi q}^{(3)} \right]_{1j} - 2 \ \left[V \hat{\alpha}_{lq}^{(3)} \right]_{\ell\ell 1j} ,\\ V_{1j} \cdot \epsilon_R^{j} &= - [\hat{\alpha}_{\varphi \varphi}]_{1j} ,\\ V_{1j} \cdot \epsilon_{s_L}^{j\ell} &= - [\hat{\alpha}_{lq}]_{\ell\ell j1}^* ,\\ V_{1j} \cdot \epsilon_{s_R}^{j\ell} &= - \left[V \hat{\alpha}_{qde}^{\dagger} \right]_{\ell\ell 1j} ,\\ V_{1j} \cdot \epsilon_{T}^{j\ell} &= - \left[\hat{\alpha}_{lq}^{\dagger} \right]_{\ell\ell j1}^* , \qquad \hat{\alpha} = \alpha \frac{v^2}{\Lambda^2} \end{split}$$

е

u

$$\begin{bmatrix} Low-E \ EFT \end{bmatrix}_{\mu=M_Z} SMEFT$$

[Cirigliano, MGA, Jenkins'2010; Cirigliano, MGA, Graesser'2012]



 $O_{lq}^{t} = (\bar{l}_{a}\sigma^{\mu\nu}e)\epsilon^{ab}(\bar{q}_{b}\sigma_{\mu\nu}u) + \text{h.c.}$

Matching with high-E EFT

$$\begin{split} \frac{\delta G_F}{G_F} &= 2 \ [\hat{\alpha}_{\varphi l}^{(3)}]_{11+22} - [\hat{\alpha}_{ll}^{(1)}]_{1221} - 2[\hat{\alpha}_{ll}^{(3)}]_{1122-\frac{1}{2}(1221)} \ , \\ V_{1j} \cdot \epsilon_L^{j\ell} &= 2 \ V_{1j} \ \left[\hat{\alpha}_{\varphi l}^{(3)} \right]_{\ell\ell} + 2 \ \left[V \hat{\alpha}_{\varphi q}^{(3)} \right]_{1j} - 2 \ \left[V \hat{\alpha}_{lq}^{(3)} \right]_{\ell\ell 1j} , \\ V_{1j} \cdot \epsilon_R^{j} &= - [\hat{\alpha}_{\varphi \varphi}]_{1j} \ , \\ V_{1j} \cdot \epsilon_{s_R}^{j\ell} &= - [\hat{\alpha}_{lq}]_{\ell\ell j1}^* \ , \\ V_{1j} \cdot \epsilon_{s_R}^{j\ell} &= - \left[V \hat{\alpha}_{qde}^{\dagger} \right]_{\ell\ell 1j} , \\ V_{1j} \cdot \epsilon_{T}^{j\ell} &= - \left[\hat{\alpha}_{lq}^{t} \right]_{\ell\ell j1}^* \ , \\ \end{split}$$

$$\begin{bmatrix} Low-E \ EFT \end{bmatrix}_{\mu=M_Z} SMEFT$$

[Cirigliano, MGA, Jenkins'2010; Cirigliano, MGA, Graesser'2012]

Beta decays sensitive to a few EFT coefficients





Scalar & tensor interactions: b_{Fierz} vs LHC





Models:

- Tree: RPV-MSSM;
- -Loop: RPC-MSSM;

[Herczeg (2001), Profumo et al (2007), Yamanaka et al. (2010)]

Scalar & tensor interactions: b_{Fierz} vs LHC





Models:

- Tree: RPV-MSSM;
- -Loop: RPC-MSSM;

[Herczeg (2001), Profumo et al (2007), Yamanaka et al. (2010)] But... Extremely hard to avoid $\pi \rightarrow lv$

- Tree: chiral theories... $(1\pm\gamma_5)$
- Loop: QED & EW mixing (S,T→P)

$$|\mathcal{A}(\pi \to \ell \nu)|^2 \sim m_\ell^2 \left(1 + \frac{M_{QCD}}{m_\ell} \epsilon_P\right)^2$$

Scalar & tensor interactions: b_{Fierz} vs LHC









[Bhattacharya et al, PRD85 (2012)]

$$N_{pp \to evX}\left(m_T^2 > m_{T,cut}^2\right) = \varepsilon \times L \times \sigma_{pp \to evX}\left(m_T^2 > m_{T,cut}^2\right) = \varepsilon \times L \times \left(\sigma_W + \sigma_S \varepsilon_S^2 + \sigma_T \varepsilon_T^2\right)$$

(Interference w/ SM ~ m/E)









$$N_{pp \to evX}\left(m_T^2 > m_{T,cut}^2\right) = \varepsilon \times \mathbf{L} \times \sigma_{pp \to evX}\left(m_T^2 > m_{T,cut}^2\right) = \varepsilon \times \mathbf{L} \times \left(\sigma_W + \sigma_S \varepsilon_S^2 + \sigma_T \varepsilon_T^2\right)$$

(Interference w/ SM ~ m/E)







$$N_{pp \to evX}\left(m_T^2 > m_{T,cut}^2\right) = \varepsilon \times L \times \sigma_{pp \to evX}\left(m_T^2 > m_{T,cut}^2\right) = \varepsilon \times L \times \left(\sigma_W + \sigma_S \varepsilon_S^2 + \sigma_T \varepsilon_T^2\right)$$

(Interference w/ SM ~ m/E)



>10⁷ 910⁶

ຊ10⁵

/s10⁴ 10³ 10²

10

10⁻¹ 10⁻²

10⁻³





$$N_{pp \to evX}\left(m_T^2 > m_{T,cut}^2\right) = \varepsilon \times \mathbf{L} \times \sigma_{pp \to evX}\left(m_T^2 > m_{T,cut}^2\right) = \varepsilon \times \mathbf{L} \times \left(\sigma_W + \sigma_S \varepsilon_S^2 + \sigma_T \varepsilon_T^2\right)$$

(Interference w/ SM \sim m/E)



Conclusions

- (Sub) permil-level precision in β decays
 - Great QCD progress
 - Experimental progress too
 - Rad. Corrections?

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- General EFT analysis available
 - \rightarrow Comparison between β -decay observables;
 - \rightarrow Comparison with APV, LEP, LHC, ...
 - $\rightarrow \beta$ decays are competitive TeV probes;

• More (Exp + Th) results expected in the near future









Backup slides







If we see a bump...

• EFT breaks down... Toy model: scalar resonance:

$$\mathcal{L} = \lambda_S V_{ud} \phi^+ \overline{u} d + \lambda_l \phi^- \overline{e} P_L \nu_e$$

• Then we have a lower-limit value for ε_{s} :

$$\sigma \cdot \mathrm{BR} \leq \frac{|V_{ud}|}{12v^2} \frac{\pi}{\sqrt{2N_c}} |\epsilon_S| \tau L(\tau)$$







$$L(\tau) = \int_{\tau}^{1} dx f_q(x) f'_q(\tau/x) / x$$

$$\tau = m^2 / s$$

$$\epsilon_S = 2\lambda_S \lambda_l \frac{v^2}{m^2}$$

Nice interplay of two experiments separated for so many orders of magnitudes!!!!

[T. Battacharya et al., 2012]

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